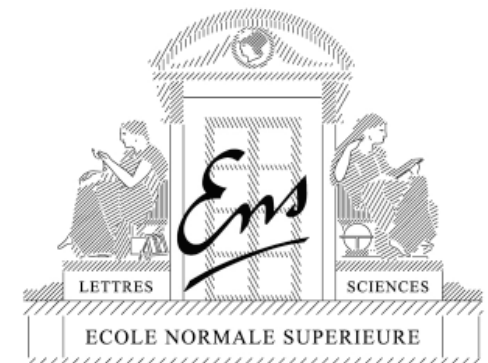


Scalars in Conformal Field Theories

Slava Rychkov



SCALARS 2011

August 26-29, 2011
Warsaw, Poland

Many reasons to think about CFT's

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UV structure Beyond the Standard Model

- ad hoc CFT hidden sectors (unparticles)
- EWSB sector (Walking and Conformal Technicolor)

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Critical phenomena

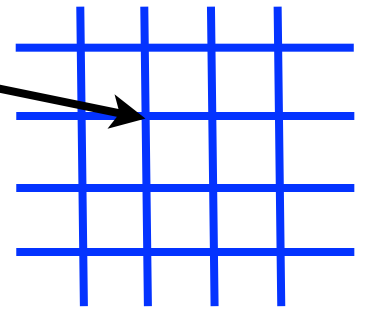
- statistical mechanics
- quantum condensed matter (=quantum criticality)

This talk

Classical $O(N)$ -invariant ferromagnet in $d=3$

$$Z = \exp \left[-\frac{1}{T} \sum_{\langle ij \rangle} \sigma_i \cdot \sigma_j \right]$$

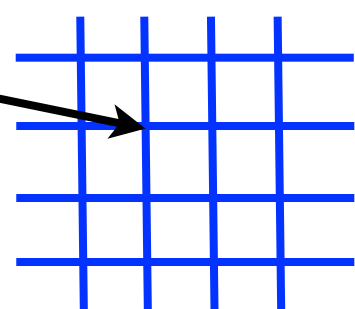
$$\sigma_i \in S^{N-1}$$



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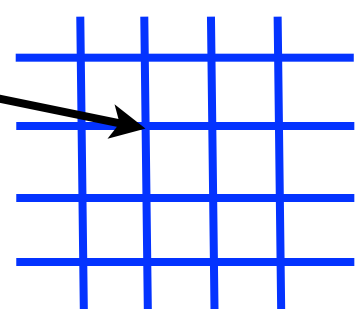
$N=1$: Ising model ($\sigma = \pm 1$)

focus on this

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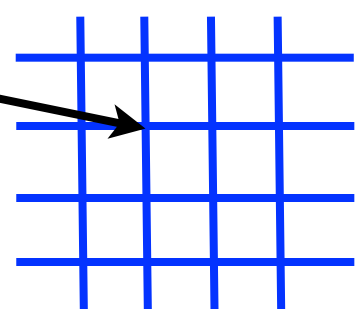
Near critical temperature:

$$\langle \sigma(r) \sigma(0) \rangle \sim \frac{1}{r^{2\Delta_\sigma}} e^{-r/\xi(T)}$$

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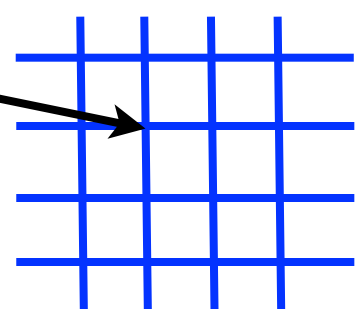
correlation length

$$\xi(T) \rightarrow \infty \quad (T \rightarrow T_c)$$

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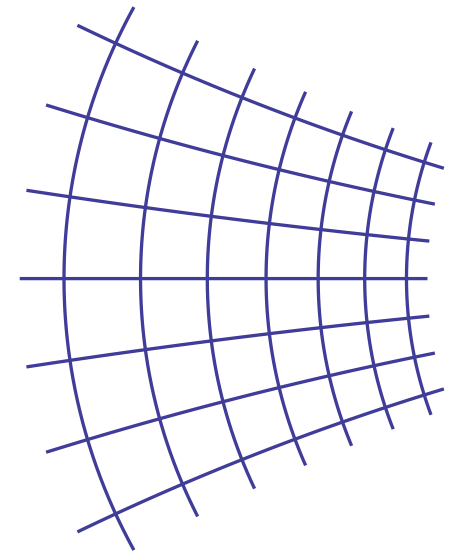
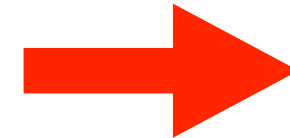
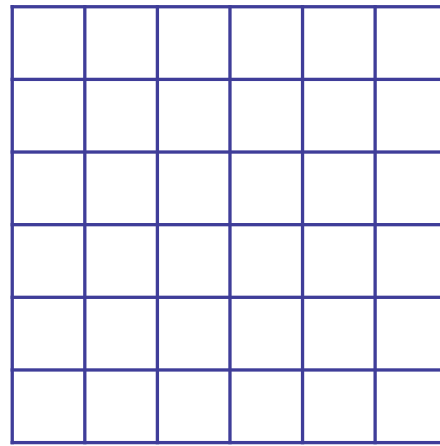
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$$\xi(T) \rightarrow \infty \quad (T \rightarrow T_c)$$

At $T=T_c$ theory is 1) scale invariant
2) conformal

Special Conformal transformation

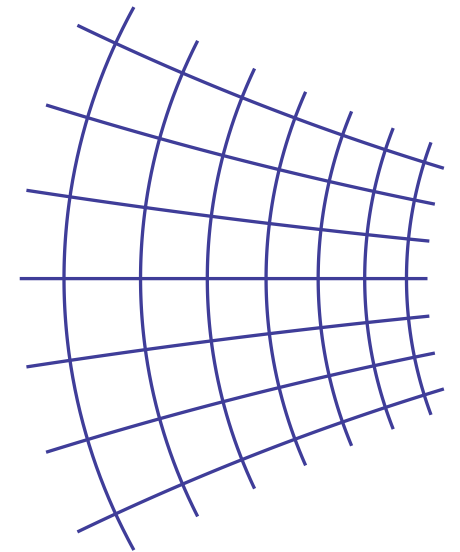
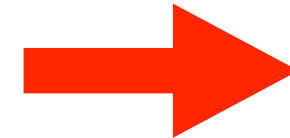
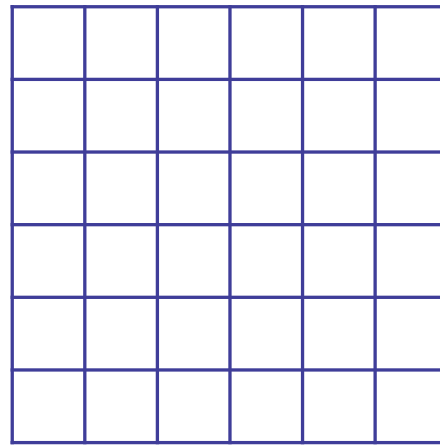
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(preserves orthogonality of coordinate grid; locally looks like dilation)

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Are there interesting scale invariant theories
without full conformal invariance?

Polchinski 1988, Dorigoni, S.R. 2009,
El-Showk, Nakayama, S.R. 2011, Antoniadis, Buican 2011
Fortin, Grinstein, Stergiou 2011

Critical Exponents (universal)

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I) Spin-spin critical exponent Δ_σ (dimension of σ) $\langle \sigma(r)\sigma(0) \rangle \sim \frac{1}{r^{2\Delta_\sigma}}$

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Related to the dimension of another local field, $\epsilon(x)$:

$$\nu = \frac{1}{d - \Delta_\epsilon}$$

Energy density field, $\varepsilon(\mathbf{x})$

$$\varepsilon \sim \sigma^2$$

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Lagrangian describing the near-critical system:

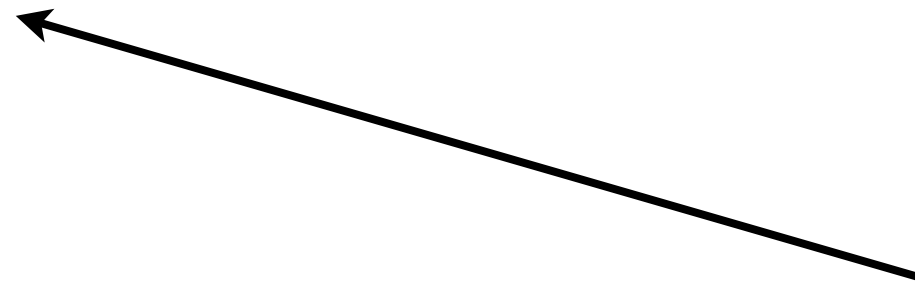
$$\mathcal{L}_{CFT} + \frac{T/T_c - 1}{a^{d-\Delta_\epsilon}} \epsilon(x)$$

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Correlation length develops at the length scale where
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$$\sim \frac{1}{\xi(T)^{d-\Delta_\epsilon}}$$

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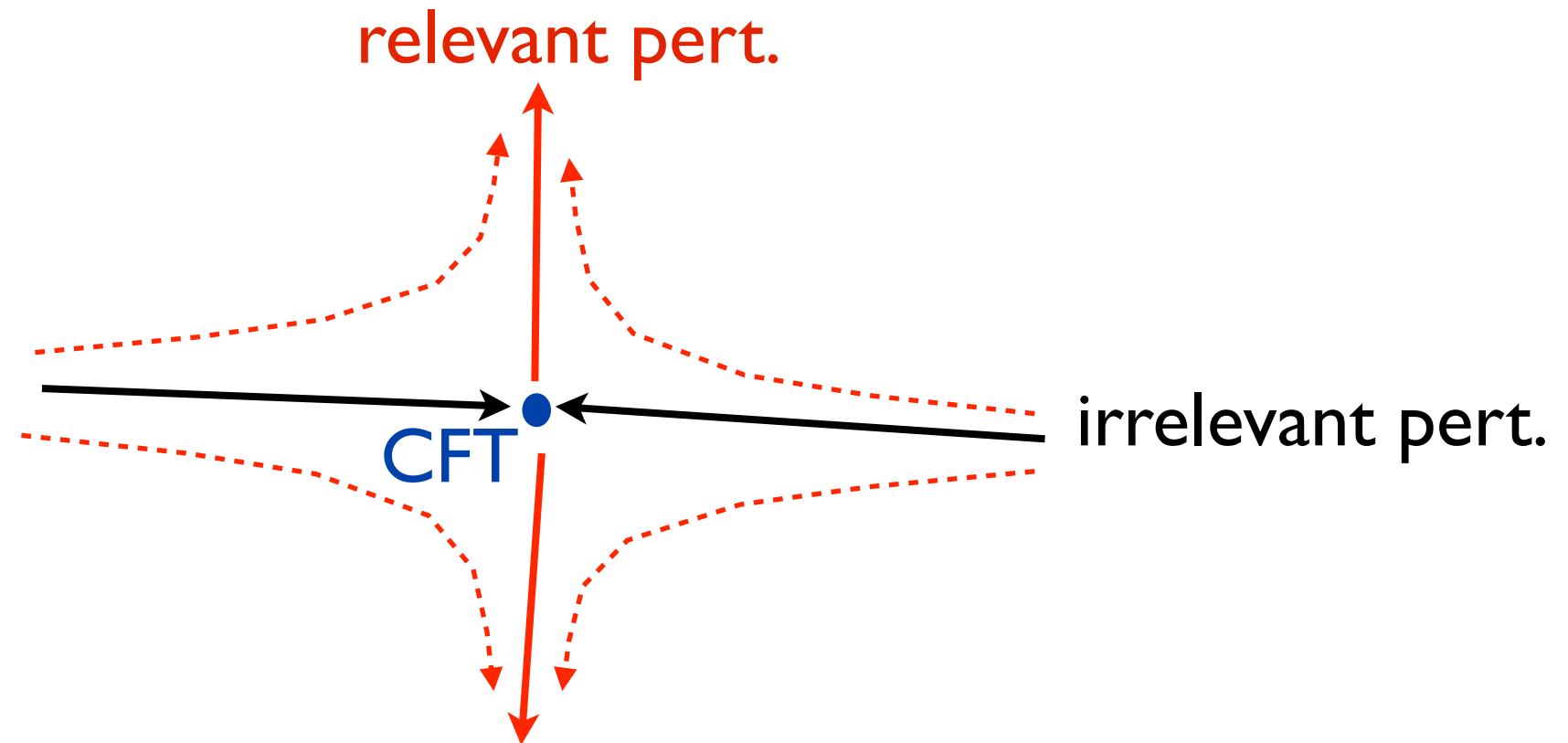
$$\Rightarrow \nu = \frac{1}{d - \Delta_\epsilon}$$

Flows

In particular from $\nu > 0$ it follows $\Delta_\varepsilon < d$ (relevant operator)

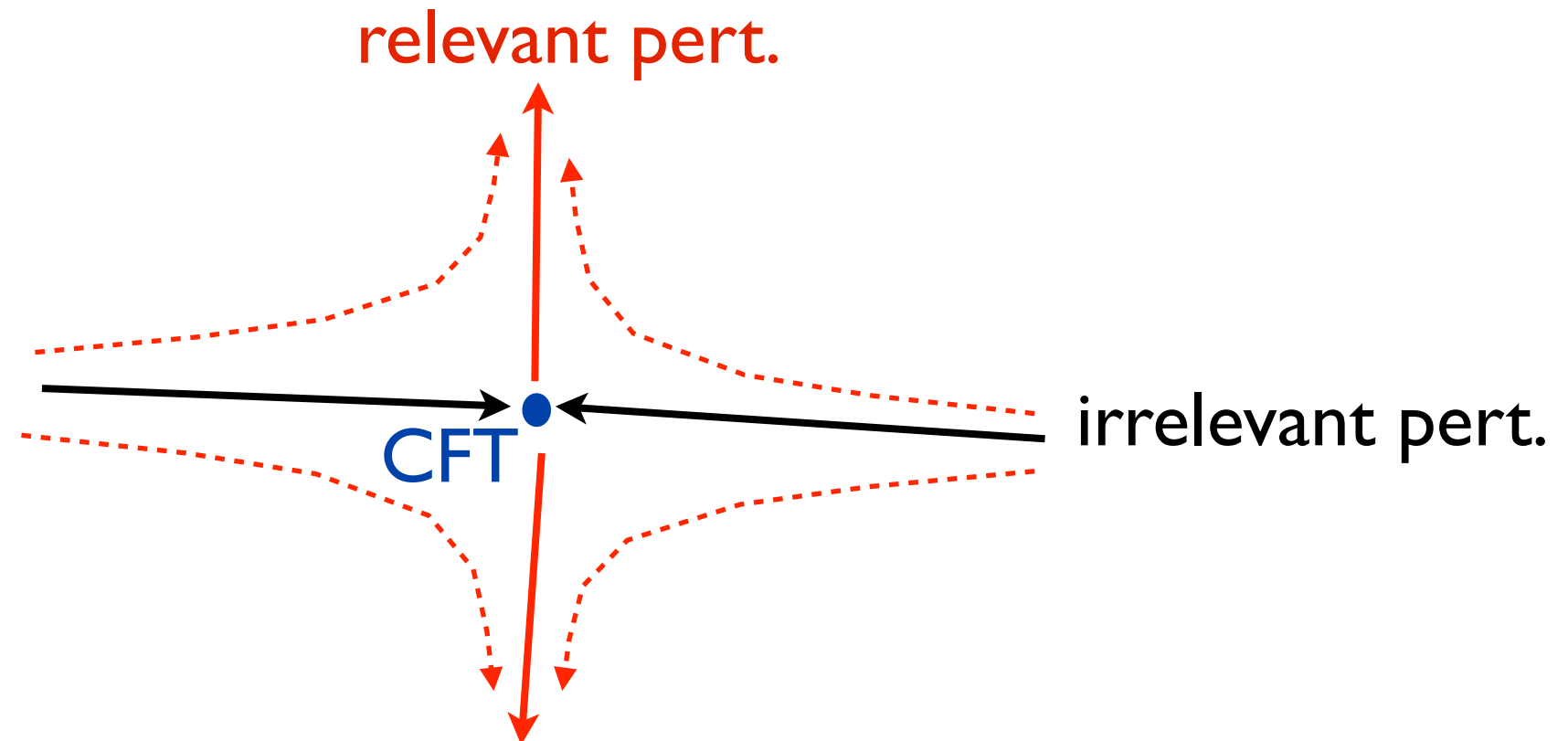
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Important: ε is **the only** relevant operator which is singlet under Z_2 symmetry $\sigma \rightarrow -\sigma$ (otherwise multicriticality)

How to determine critical exponents

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$$\Delta_{\varepsilon} = 1.412(1)$$

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- Field theory techniques in $4-\varepsilon$ dimensions

ϵ -expansion Wilson, Fischer

Study scalar field theory in $4-\epsilon$ dimensions

$$\mathcal{L} = (\partial\phi)^2 + \lambda\phi^4$$

$$\beta_\lambda = -\epsilon\lambda + \frac{\lambda^2}{16\pi^2} + \dots \rightarrow 0$$

$$\frac{\lambda_*}{16\pi^2} = O(\epsilon) \ll 1 \quad \text{weakly coupled fixed point}$$

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Compute critical exponents at $\epsilon \ll 1$ and then extrapolate to $\epsilon = 1$

Works pretty well but not to arbitrary accuracy (divergent series)

I think the epsilon expansion ended the subject in the practical sense.

You can calculate more or less what you want with good accuracy but aesthetically the subject is not closed yet.

It's possible that there will be classification of fixed points in three dimensions. But that's just dreams.

A.M. Polyakov, 2003 interview



Conformal bootstrap

Ferrara, Gatto, Grillo 1973

Polyakov 1974

Consistency eq. four 4pt
function

$$\langle \sigma(x_1) \sigma(x_2) \sigma(x_3) \sigma(x_4) \rangle$$

Conformal bootstrap

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OPE:

$$\sigma(x_1) \sigma(x_2) = \mathbf{1} + \epsilon + \dots$$

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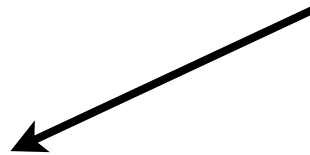
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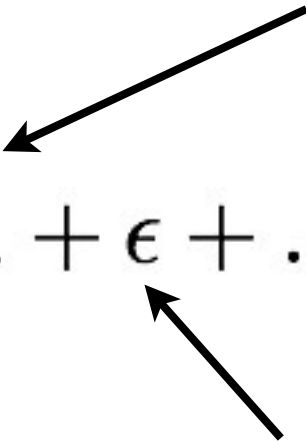
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The diagram shows two arrows pointing towards the OPE expansion. One arrow originates from the expression $\langle \sigma \sigma \rangle \neq 0$ and points to the identity term $\mathbf{1}$ in the expansion. The other arrow originates from the expression $\langle \sigma \sigma \epsilon \rangle \neq 0$ and points to the ϵ term in the expansion.

$$\langle \sigma \sigma \rangle \neq 0$$

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$$\langle \sigma \sigma \epsilon \rangle \neq 0$$

Conformal symmetry fixes:

$$\langle \sigma(x) \sigma(y) \epsilon(0) \rangle = \frac{\lambda}{|x - y|^{2\Delta_\sigma - \Delta_\epsilon} |x|^{2\Delta_\epsilon} |y|^{2\Delta_\epsilon}}$$

Conformal OPE

$$\sigma(x_1) \sigma(x_2) = \sum_O \lambda_O C(x_1 - x_2, \partial_{x_2}) O(x_2)$$

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fixed by conformal symmetry



$$O = O_{\Delta}^{(l)}$$

$$l = 2, 4, 6, \dots$$

$$\Delta \geq l + d + 2 \quad (\geq d/2 - 1, l = 0)$$

$$\langle \sigma(x_1) \sigma(x_2) \sigma(x_3) \sigma(x_4) \rangle$$



OPE



OPE

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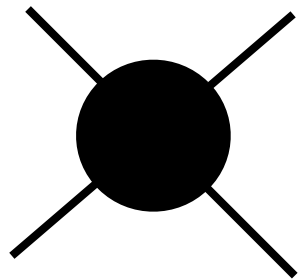


OPE



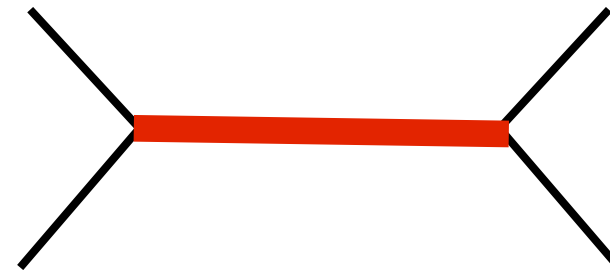
OPE

Conformal partial wave



=

$$\sum_o \lambda_o^2$$



$$\langle \sigma(x_1) \sigma(x_2) \sigma(x_3) \sigma(x_4) \rangle$$

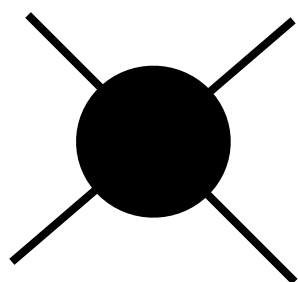


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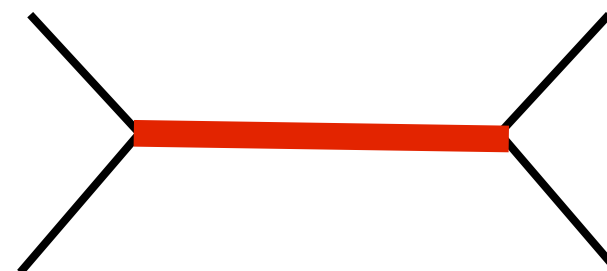
OPE

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$$\langle \sigma \sigma \sigma \sigma \rangle = \frac{g(u, v)}{r_{12}^{2\Delta_\sigma} r_{34}^{2\Delta_\sigma}}$$

$$u = \left(\frac{r_{12} r_{34}}{r_{13} r_{24}} \right)^2 \quad v = \left(\frac{r_{14} r_{23}}{r_{13} r_{24}} \right)^2$$

$$\langle \sigma(x_1) \sigma(x_2) \sigma(x_3) \sigma(x_4) \rangle$$

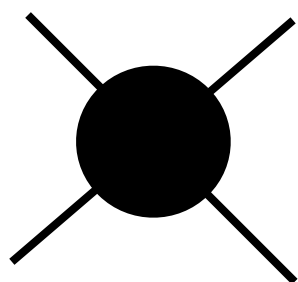


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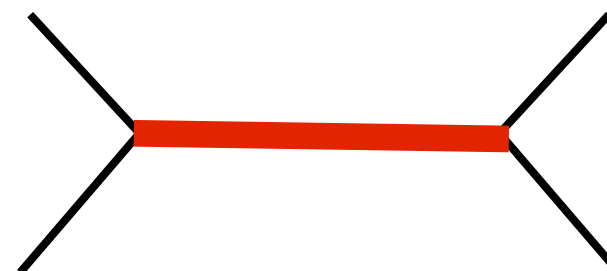
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$$g(u, v) = 1 + \sum_O \lambda_O^2 g_O(u, v)$$



known functions of u, v

Crossing symmetry/OPE associativity

$$\begin{array}{ccc} \langle \sigma(x_1) \sigma(x_2) \sigma(x_3) \sigma(x_4) \rangle & = & \langle \sigma(x_1) \sigma(x_2) \sigma(x_3) \sigma(x_4) \rangle \\ \underbrace{\hspace{1.5cm}}_{\text{OPE}} \quad \underbrace{\hspace{1.5cm}}_{\text{OPE}} & & \underbrace{\hspace{3cm}}_{\text{OPE}} \\ & & \underbrace{\hspace{1.5cm}}_{\text{OPE}} \end{array}$$

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Conformal bootstrap equation for CFT couplings and spectrum
(no progress for 30 years)

Resurrecting Conformal Bootstrap

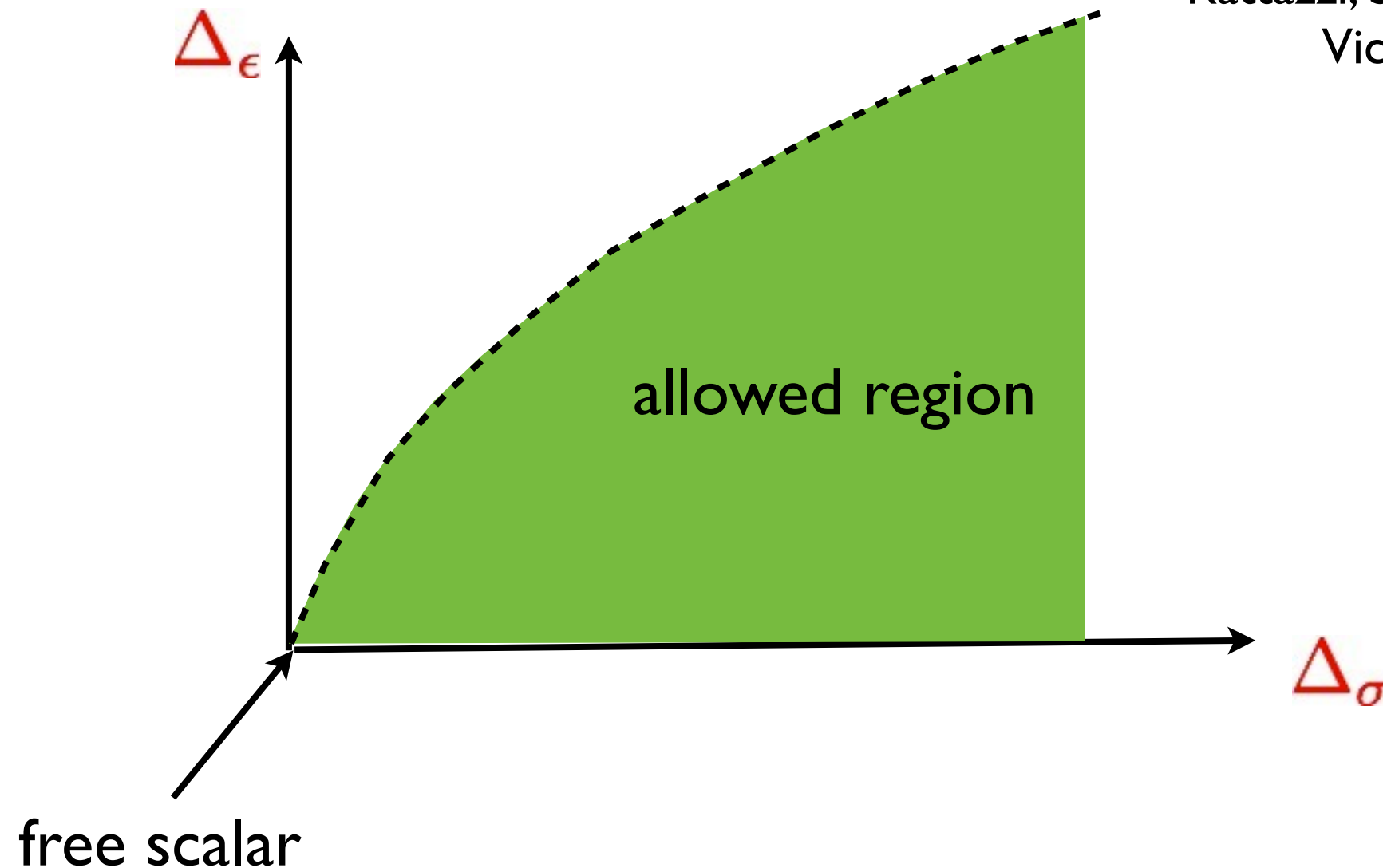
Rattazzi, S.R., Tonni, Vichi 2008

S.R., Vichi 2009

Caracciolo, S.R. 2009

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Vichi 2011



Resurrecting Conformal Bootstrap

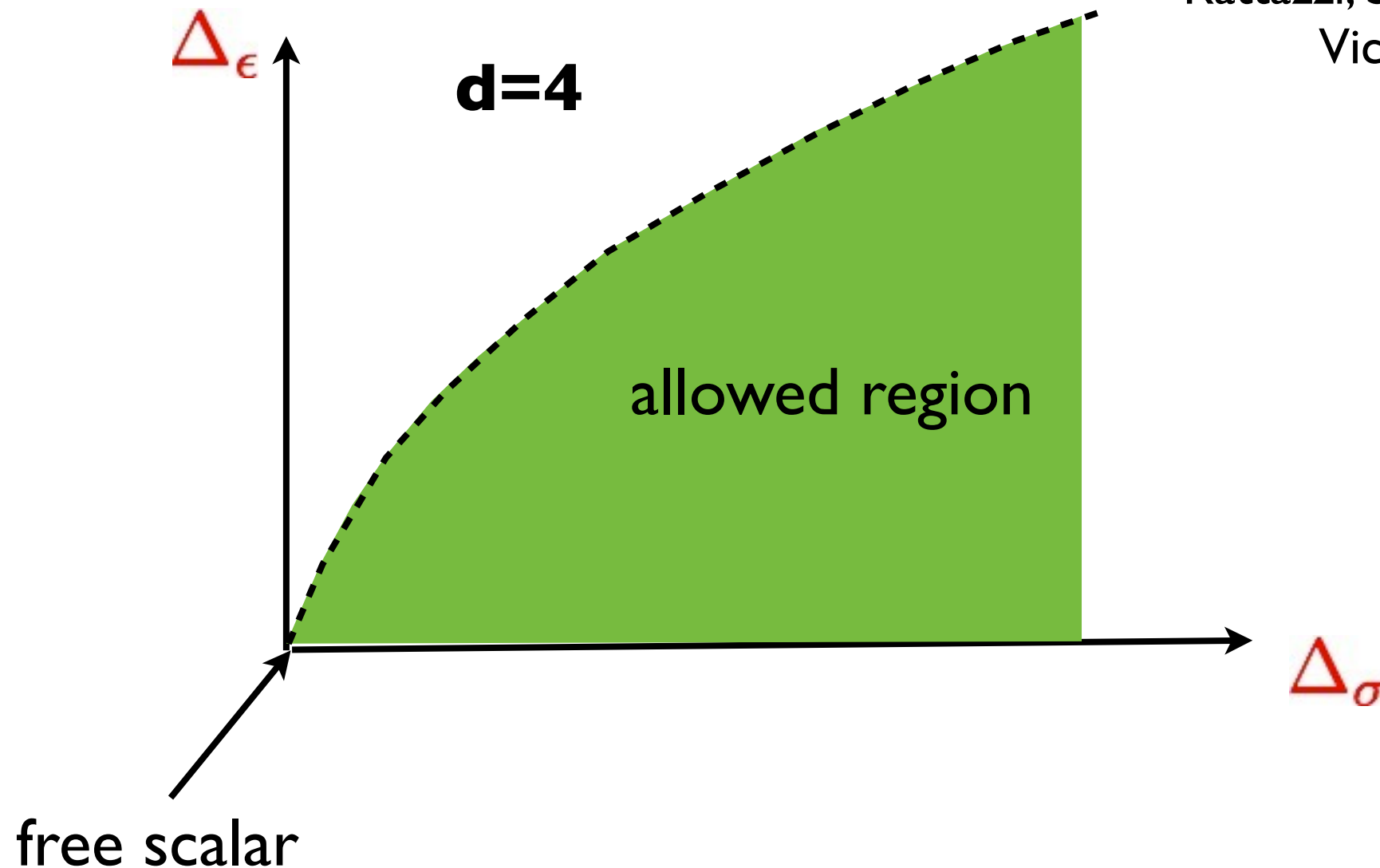
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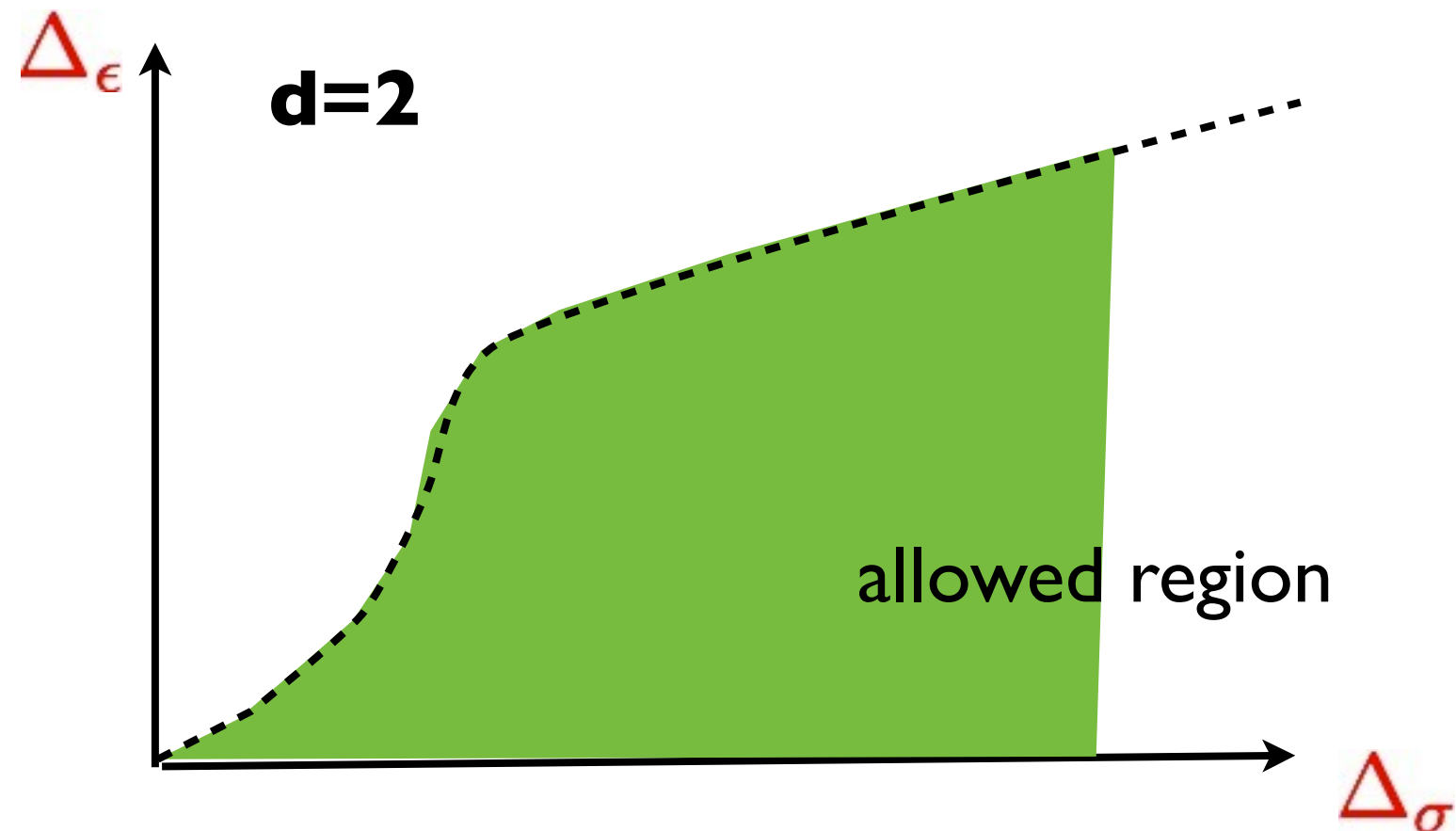
Vichi 2011



Motivated by Conformal Technicolor

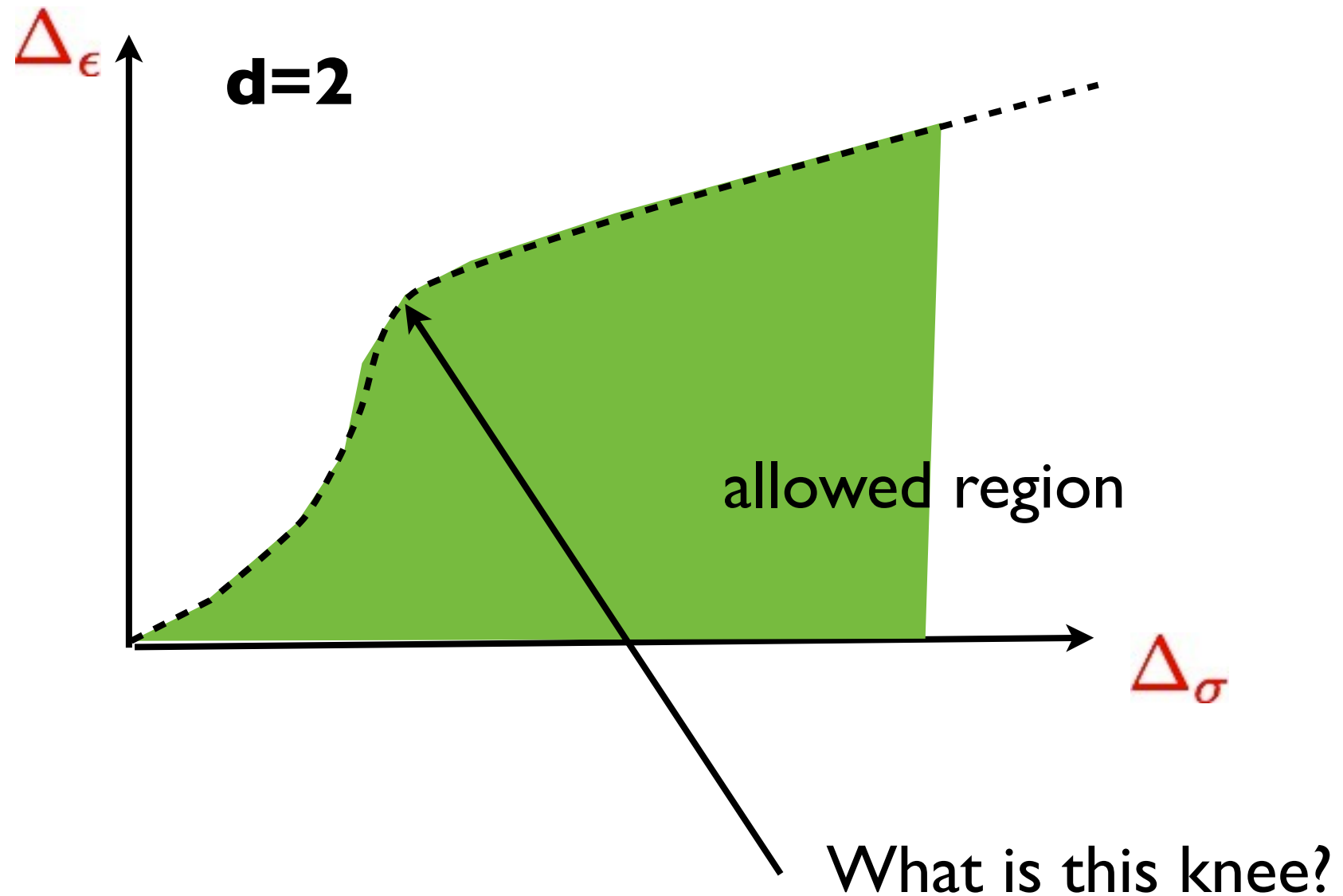
Exercise in $d=2$

S.R., Vichi 2009



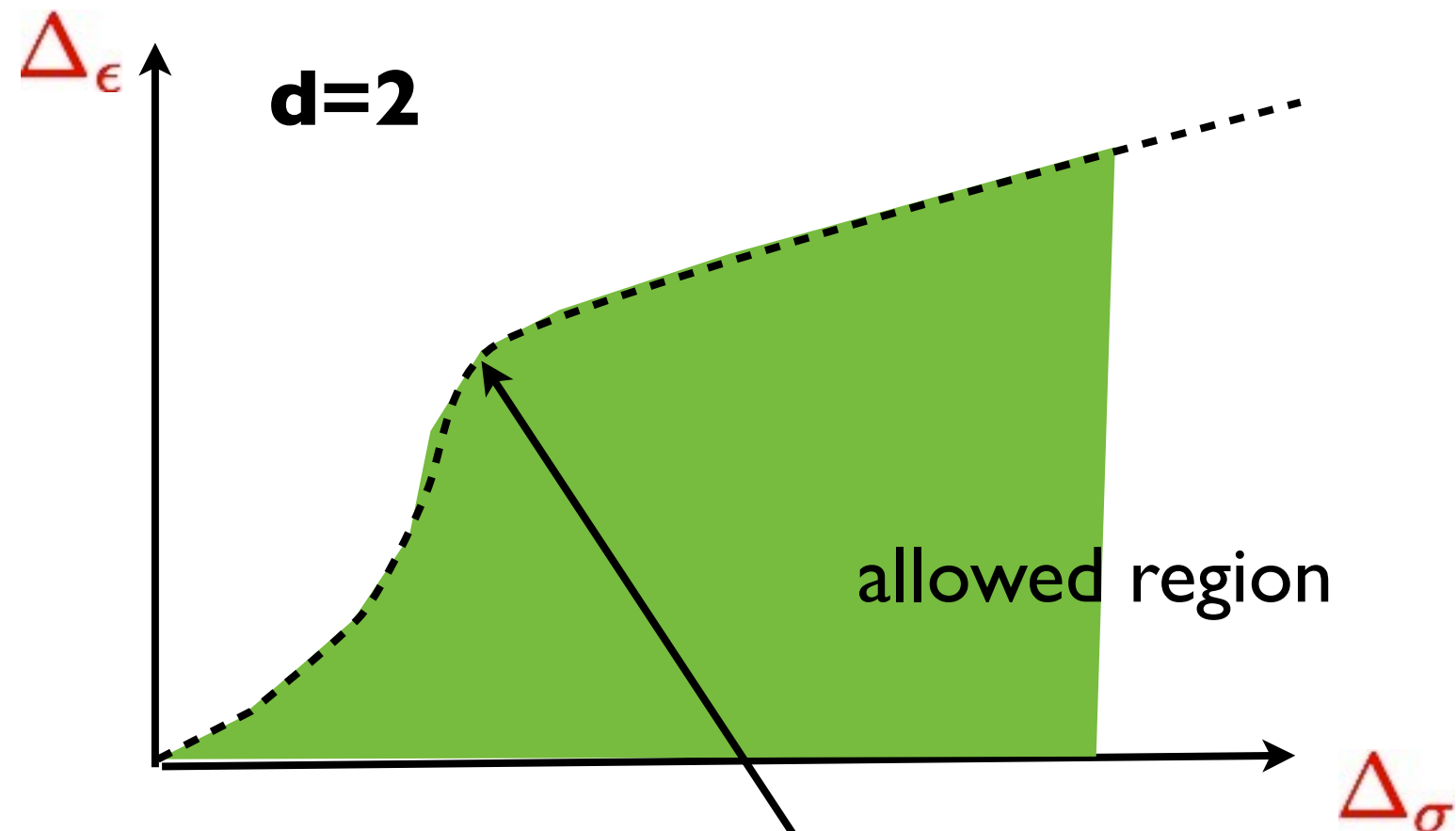
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Exercise in $d=2$

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What is this knee?

It's the 2D Ising model!

$$\Delta_\sigma = 1/8, \Delta_\epsilon = 1$$

Extracting d=3 Ising critical exponents

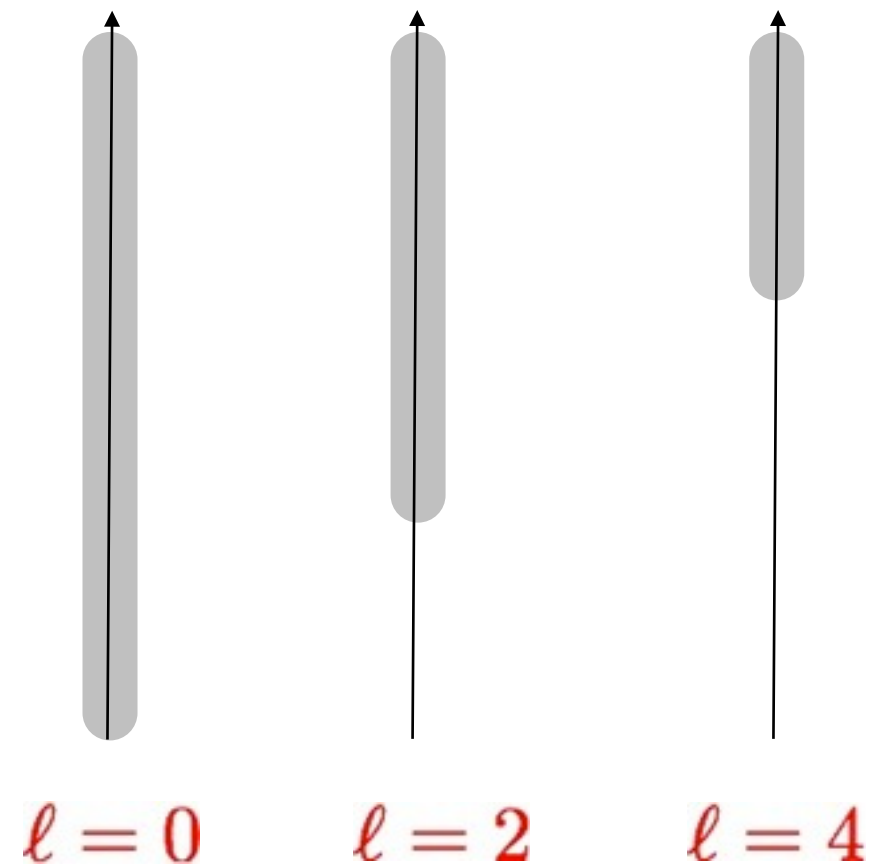
Idea 0:

Look for the knee on the boundary of allowed region in $(\Delta_\sigma, \Delta_\varepsilon)$ plane

Origin of the knee

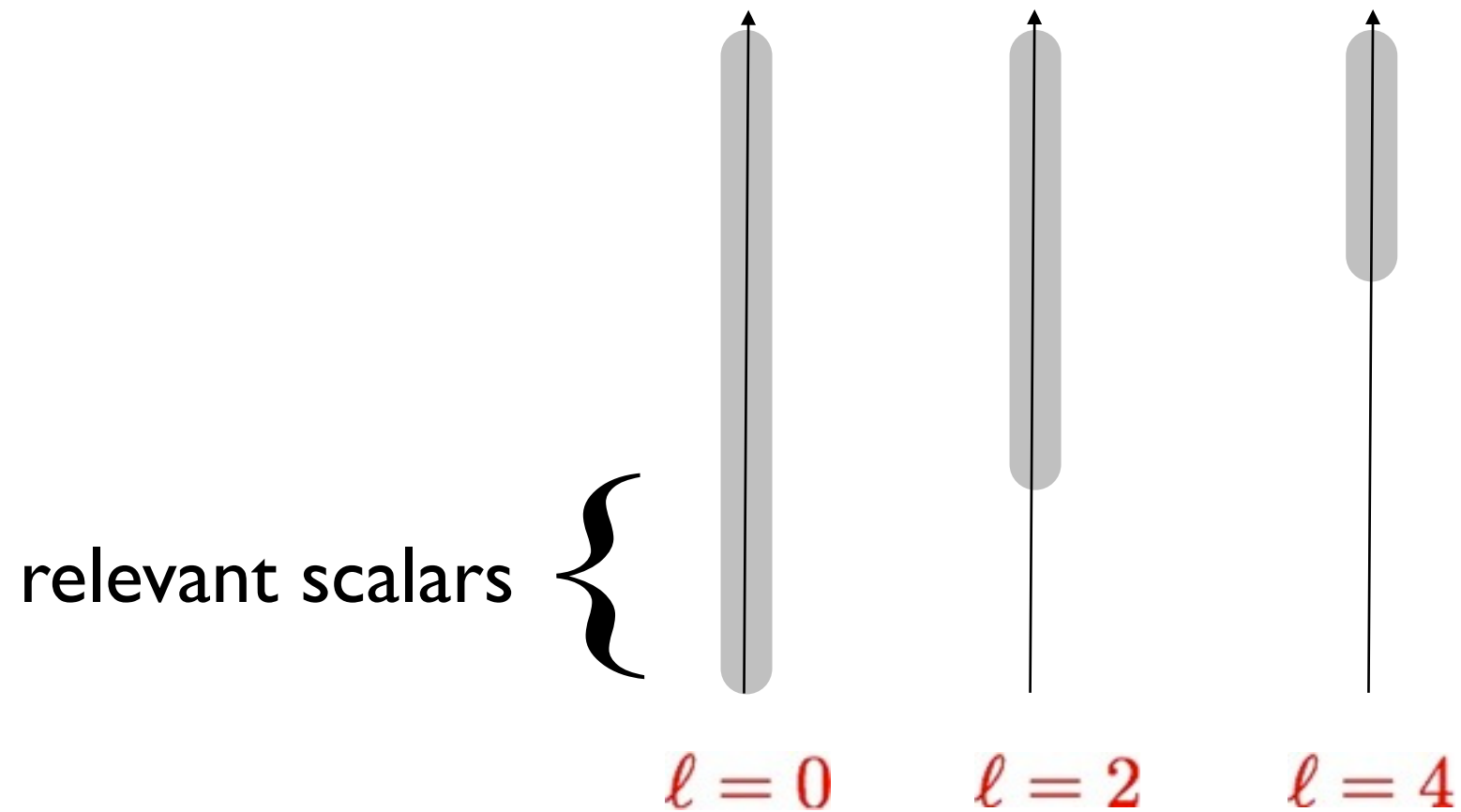
Origin of the knee

Allowed OPE spectrum:



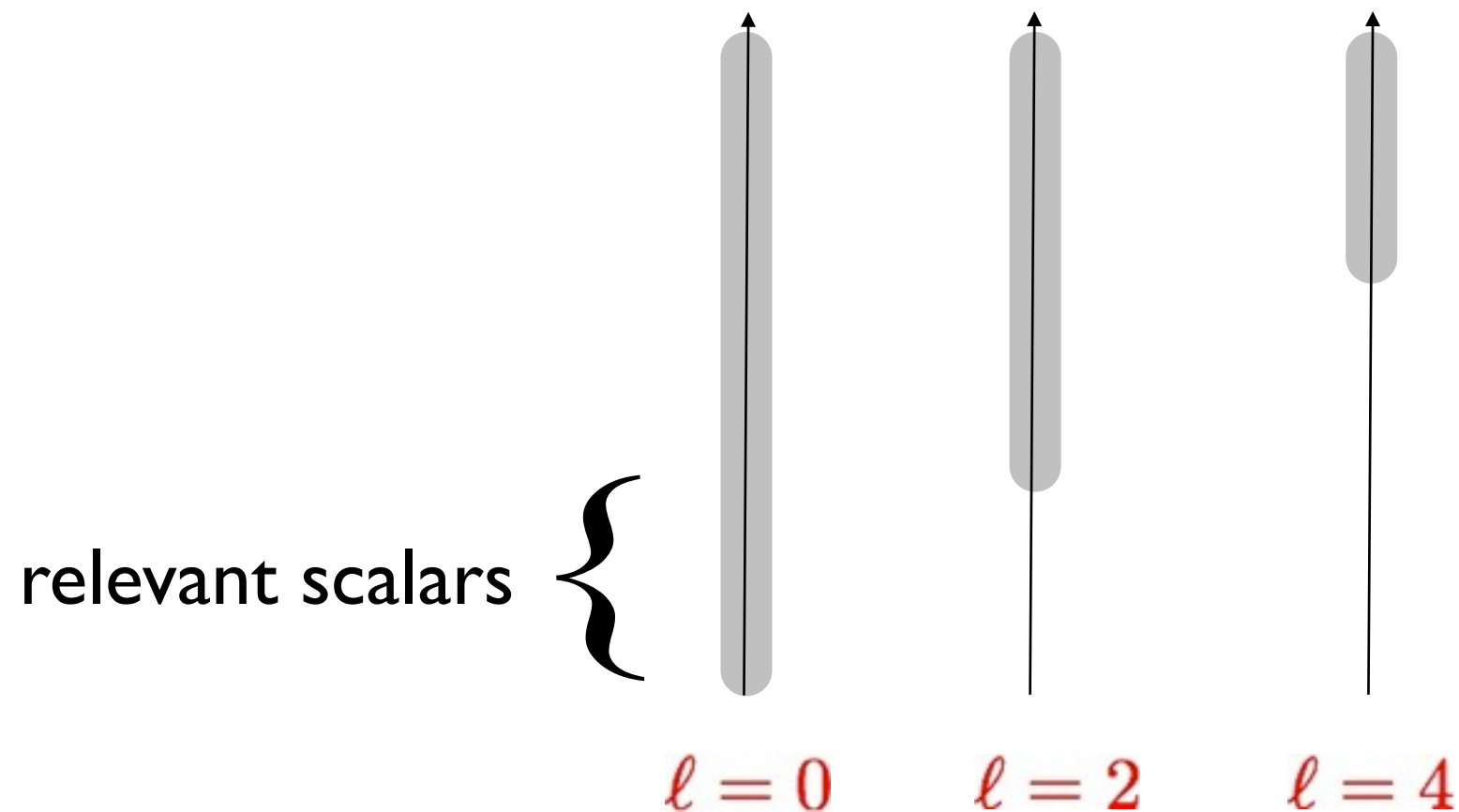
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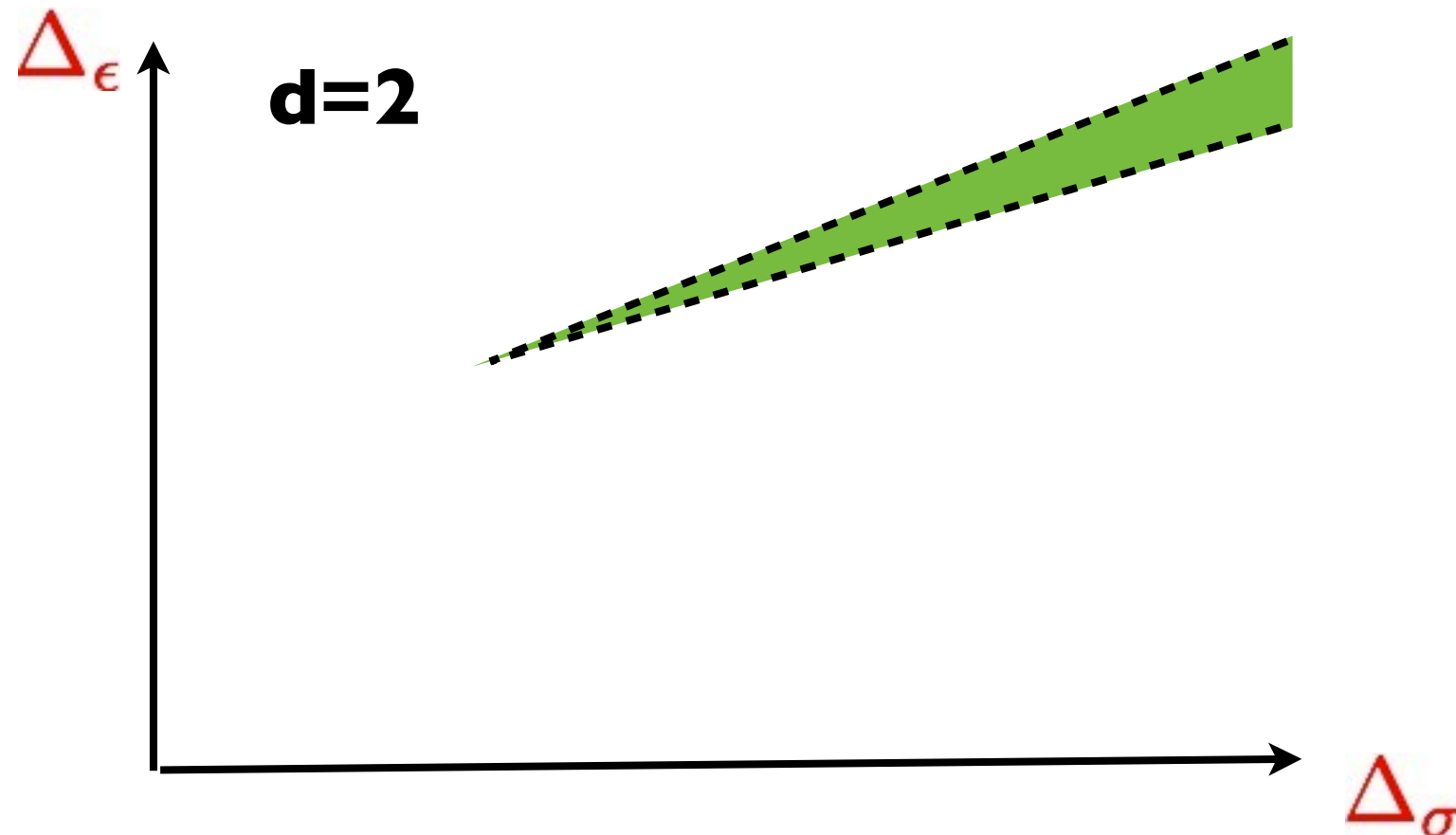
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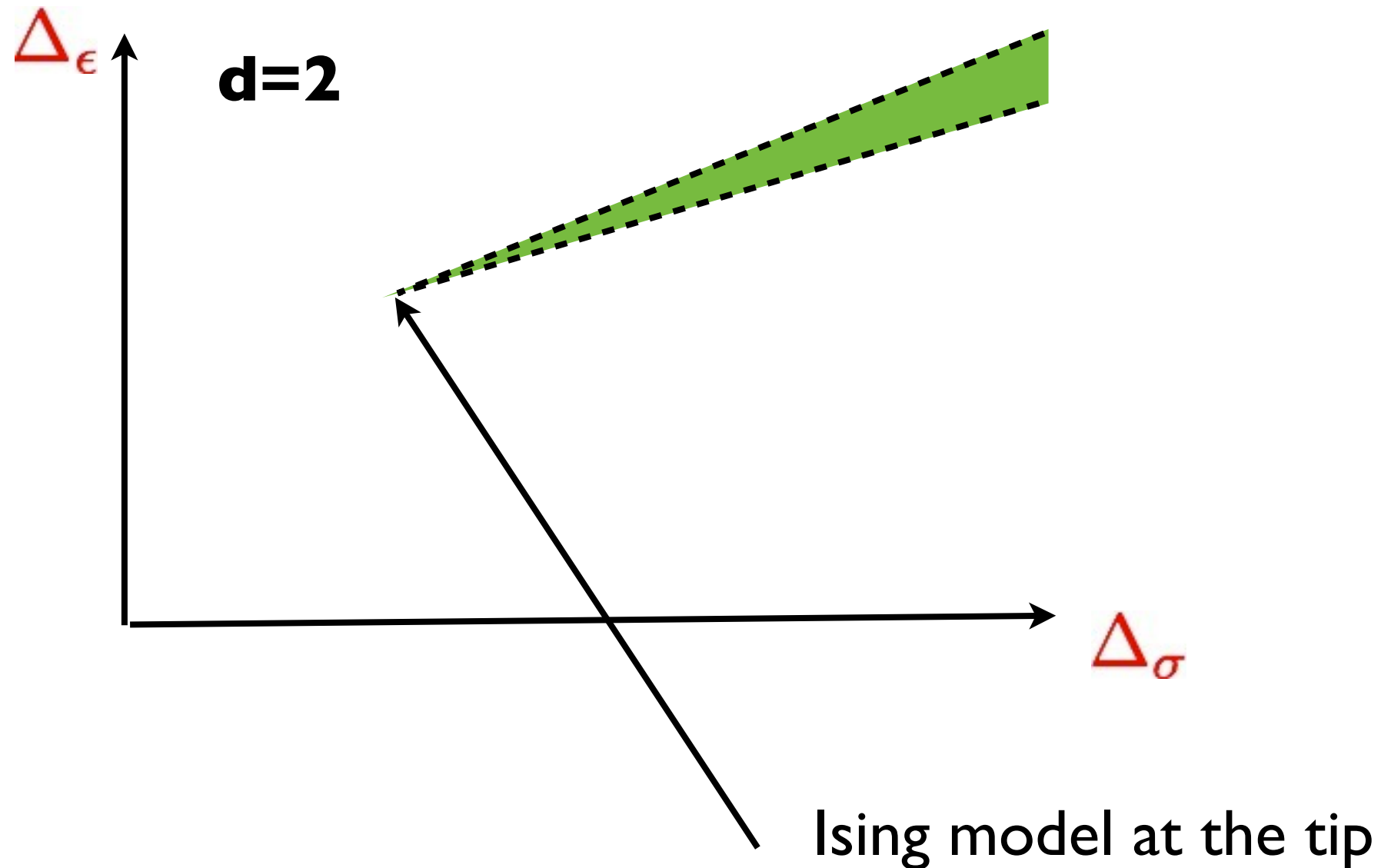
For $\Delta_\sigma < 1/8$ solutions to crossing symmetry have always ≥ 2
relevant scalars
(while Ising model has only one)

Impose condition that no other scalars in OPE
(apart from ε) with $\Delta < 2.5$

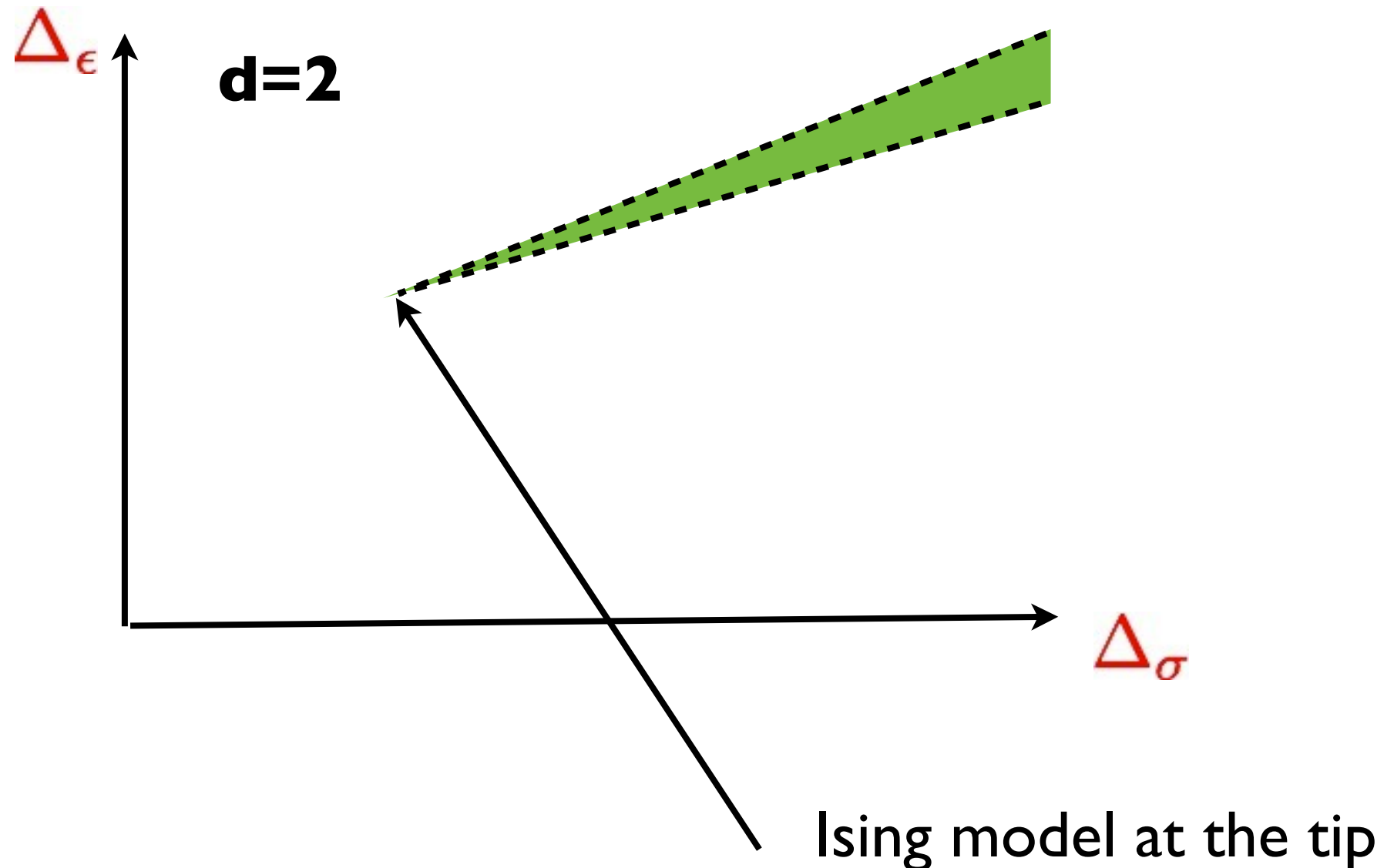
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Impose condition that no other scalars in OPE
(apart from ϵ) with $\Delta < 2.5$



Allows much sharper determination of critical exponents

Why not yet applied in d=3?

Explicit conformal partial waves: Dolan, Osborn 2001

$$\mathbf{d=4} \quad g_O(u, v) = \frac{z\bar{z}}{z - \bar{z}} [k_{\Delta+l}(z)k_{\Delta-l-2}(\bar{z}) - (z \leftrightarrow \bar{z})]$$

$$\mathbf{d=2} \quad g_O(u, v) = k_{\Delta+l}(z)k_{\Delta-l}(\bar{z}) + (z \leftrightarrow \bar{z})$$

$$u = z\bar{z}, \quad v = (1 - z)(1 - \bar{z})$$

$$k_\beta(x) \equiv x^{\beta/2} {}_2F_1\left(\frac{\beta}{2}, \frac{\beta}{2}, \beta; x\right)$$

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$$\mathbf{d=2} \quad g_O(u, v) = k_{\Delta+l}(z)k_{\Delta-l}(\bar{z}) + (z \leftrightarrow \bar{z})$$

$$u = z\bar{z}, \quad v = (1 - z)(1 - \bar{z})$$

$$k_\beta(x) \equiv x^{\beta/2} {}_2F_1\left(\frac{\beta}{2}, \frac{\beta}{2}, \beta; x\right)$$

In $\mathbf{d=3}$ equally simple expressions are not yet known

There exist double power series in $(u, 1-v)$ which can be used but with more difficulty