Vector bosons scattering and boundary conditions in KK toy model

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Standard Model Higgs mechanism, tree-level unitarity





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Higher-dimensional models Compact extra dimension(s)

• Infinite space dimension compactified e.g. on S^1/Z_2 orbifold.



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• Finite space dimension - "interval approach".

- Gauge Higgs Unification Higgs field arises from extra component(s) of higher-dimensional gauge field.
- Higgsless models Electroweak symmetry is broken via a non-trivial choice of boundary conditions for gauge field, no physical scalar is needed.

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Pure gauge 5D theory on flat background with one extra finite space dimension

- $y \in (0, \pi R)$
- $A^a_N(x^\mu, y) = (A^a_\nu(x^\mu, y), A^a_5(x^\mu, y))$
- $F^a_{MN} = \partial_M A^a_N \partial_N A^a_M + g_5 f^{abc} A^b_M A^c_N$

•
$$\mathscr{L}_{gauge} = -\frac{1}{4}F^a_{MN}F^{aMN} = -\frac{1}{4}F^a_{\mu\nu}F^{a\mu\nu} - \frac{1}{2}F^a_{\mu5}F^{a\mu5}$$

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•
$$R_{\xi}$$
 gauge: $\mathscr{L}_{\text{g.f.}} = -\frac{1}{2\xi} \left(\partial_{\mu} A^{a\mu} - \xi \partial_{5} A^{a}_{5} \right)^{2}$

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Equations of motion, boundary conditions

Variational principle of least action ($\delta S = 0$) gives us

equations of motion

•
$$\partial^M F^a_{M\nu} - g_5 f^{abc} F^b_{M\nu} A^{cM} + \frac{1}{\xi} \partial_\nu \partial^\mu A^a_\mu - \partial^\nu \partial_5 A^a_5 = 0$$

• $\partial^\mu F^a_{\mu 5} - g_5 f^{abc} F^b_{\mu 5} A^{c\mu} + \partial_5 \partial^\mu A^a_\mu - \xi \partial_5 \partial_5 A^a_5 = 0$

• requirement of zero boundary terms

•
$$[(\partial_{\mu}A^{a\mu} - \xi \partial_{5}A^{a}_{5}) \delta A^{a}_{5}]^{\pi R}_{0} = 0$$

• $[F^{a\nu 5} \delta A^{a}_{\nu}]^{\pi R}_{0} = 0$

Possible boundary conditions

•
$$A^a_{\mu}|_{y=0,\pi R} = 0$$
, $\partial_5 A^a_5|_{y=0,\pi R} = 0$
• $\partial_5 A^a_{\mu}|_{y=0,\pi R} = 0$, $A^a_5|_{y=0,\pi R} = 0$

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KK expansion, effective 4D Lagrangian

- KK expansion: $A^a_\mu(x,y) = \sum_n A^{a,n}_\mu(x) \, \varphi_{a,n}(y)$
- Boundary conditions keep the operator ∂_y^2 hermitian with respect to scalar product $\langle f,g\rangle = \int_0^{\pi R} f^*(y)g(y)\,\mathrm{d}y$ \Rightarrow complete orthonormal basis exists and satisfies $\varphi_{a,n}''(y) = -m_{a,n}^2\varphi_{a,n}(y)$

•
$$\Box_5 A^a_\nu(x,y) - \partial_\nu \partial^\mu A^a_\mu(x,y) = 0 \Leftrightarrow (\Box_4 + m^2_{a,n}) A^{a,n}_\nu(x) - \partial_\nu \partial^\mu A^{a,n}_\mu(x) = 0$$

• effective 4D Lagrangian: $\mathscr{L}_{4D}(x) = \int_0^{\pi R} \mathscr{L}(x, y) dy$

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Interactions of effective 4D fields



- In unitary gauge (ξ → +∞) are all massive modes A₅^{a,n} (n > 0) unphysical and play similar role as would-be Goldstone bosons of SM Higgs mechanism.
- Effective coupling constant in every vertex is given by an integral of wave functions, e.g. for 3V vertex:
 g_{(a,i)(b,j)(c,k)} = g₅ ∫₀^{πR} φ_{a,i} φ_{b,j} φ_{c,k} dy
- Scattering amplitudes of general $V_L V_L \rightarrow V_L V_L$ process do not grow explicitly with energy for any allowed choice of boundary conditions.

Example

Boundary conditions:

• y = 0: $A_{\nu}^{1,2} = 0$, $\partial_y A_{\nu}^3 = 0$

•
$$y = \pi R$$
: $\partial_y A^a_\nu = 0$

Wave functions and masses:

•
$$\varphi_{W,n}(y) = \sqrt{\frac{2}{\pi R}} \sin \frac{(2n-1)y}{2R} \Rightarrow m_{W,n} = \frac{(2n-1)}{2R}, \ n \ge 1$$

• $\varphi_{Z,n}(y) = \sqrt{\frac{2}{2^{\delta_{k,0}} \pi R}} \cos \frac{2ny}{2R} \Rightarrow m_{Z,n} = \frac{2n}{2R}, \ n \ge 0$

Coupling constant:

•
$$g_{(W,i)(W,j)(Z,k)} = \frac{g_5}{\sqrt{2^{\delta_{k,0}}\pi R}} \left(\delta_{j,i+k} + \delta_{j,i-k} - \delta_{j,1-i+k} \right)$$

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Final remarks

- There are several ways how to break electroweak symmetry in higher-dimensional models.
- 5D theory is non-renormalizable, unitarity breakdown is postponed to the cutoff scale of effective 4D theory (related to the size of extra dimension).

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Realistic models:

- More complicated gauge group, e.g. $SU(2)_R \times SU(2)_L \times U(1)_{B-L}$.
- Warped extra dimensions.

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Thank you for your attention.

Vector bosons scattering and boundary conditions in KK toy model arXiv:1106.3043 [hep-th]

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