# Vector bosons scattering and boundary conditions in KK toy model 

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## Standard Model

Higgs mechanism, tree-level unitarity



## Higher-dimensional models

## Compact extra dimension(s)

- Infinite space dimension compactified e.g. on $S^{1} / Z_{2}$ orbifold.

- Finite space dimension - "interval approach".


## Higher-dimensional models

Electroweak symmetry breaking

- Gauge Higgs Unification - Higgs field arises from extra component(s) of higher-dimensional gauge field.
- Higgsless models - Electroweak symmetry is broken via a non-trivial choice of boundary conditions for gauge field, no physical scalar is needed.


## Pure gauge 5D theory on flat background

 with one extra finite space dimension- $y \in(0, \pi R)$
- $A_{N}^{a}\left(x^{\mu}, y\right)=\left(A_{\nu}^{a}\left(x^{\mu}, y\right), A_{5}^{a}\left(x^{\mu}, y\right)\right)$
- $F_{M N}^{a}=\partial_{M} A_{N}^{a}-\partial_{N} A_{M}^{a}+g_{5} f^{a b c} A_{M}^{b} A_{N}^{c}$
- $R_{\xi}$ gauge: $\mathscr{L}_{\text {g.f. }}=-\frac{1}{2 \xi}\left(\partial_{\mu} A^{a \mu}-\xi \partial_{5} A_{5}^{a}\right)^{2}$


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- $F_{M N}^{a}=\partial_{M} A_{N}^{a}-\partial_{N} A_{M}^{a}+g_{5} f^{a b c} A_{M}^{b} A_{N}^{c}$
- $\mathscr{L}_{\text {gauge }}=-\frac{1}{4} F_{M N}^{a} F^{a M N}=-\frac{1}{4} F_{\mu \nu}^{a} F^{a \mu \nu}-\frac{1}{2} F_{\mu 5}^{a} F^{a \mu 5}$
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## Equations of motion, boundary conditions

Variational principle of least action $(\delta S=0)$ gives us

- equations of motion
- $\partial^{M} F_{M \nu}^{a}-g_{5} f^{a b c} F_{M \nu}^{b} A^{c M}+\frac{1}{\xi} \partial_{\nu} \partial^{\mu} A_{\mu}^{a}-\partial^{\nu} \partial_{5} A_{5}^{a}=0$
- $\partial^{\mu} F_{\mu 5}^{a}-g_{5} f^{a b c} F_{\mu 5}^{b} A^{c \mu}+\partial_{5} \partial^{\mu} A_{\mu}^{a}-\xi \partial_{5} \partial_{5} A_{5}^{a}=0$
- requirement of zero boundary terms
- $\left[\left(\partial_{\mu} A^{a \mu}-\xi \partial_{5} A_{5}^{a}\right) \delta A_{5}^{a}\right]_{0}^{\pi R}=0$
- $\left[F^{a \nu 5} \delta A_{\nu}^{a}\right]_{0}^{\pi R}=0$

Possible boundary conditions


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Possible boundary conditions

- $\left.\delta A_{\mu}^{a}\right|_{y=0, \pi R}=0,\left.\delta A_{5}^{a}\right|_{y=0, \pi R}=0$


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Possible boundary conditions

- $\left.A_{\mu}^{a}\right|_{y=0, \pi R}=0,\left.\partial_{5} A_{5}^{a}\right|_{y=0, \pi R}=0$
- $\left.\partial_{5} A_{\mu}^{a}\right|_{y=0, \pi R}=0,\left.A_{5}^{a}\right|_{y=0, \pi R}=0$


## KK expansion, effective 4D Lagrangian

- KK expansion: $A_{\mu}^{a}(x, y)=\sum_{n} A_{\mu}^{a, n}(x) \varphi_{a, n}(y)$
- Boundary conditions keep the operator $\partial_{y}^{2}$ hermitian with respect to scalar product $\langle f, g\rangle=\int_{0}^{\pi R} f^{*}(y) g(y) \mathrm{d} y$ $\Rightarrow$ complete orthonormal basis exists and satisfies

$$
\varphi_{a, n}^{\prime \prime}(y)=-m_{a, n}^{2} \varphi_{a, n}(y)
$$

- $\square_{5} A_{\nu}^{a}(x, y)-\partial_{\nu} \partial^{\mu} A_{\mu}^{a}(x, y)=0 \Leftrightarrow$ $\left(\square_{4}+m_{a, n}^{2}\right) A_{\nu}^{a, n}(x)-\partial_{\nu} \partial^{\mu} A_{\mu}^{a, n}(x)=0$
- effective 4 D Lagrangian: $\mathscr{L}_{4 \mathrm{D}}(x)=\int_{0}^{\pi R} \mathscr{L}(x, y) \mathrm{d} y$


## Interactions of effective 4D fields



- In unitary gauge $(\xi \rightarrow+\infty)$ are all massive modes $A_{5}^{a, n}$ ( $n>0$ ) unphysical and play similar role as would-be Goldstone bosons of SM Higgs mechanism.
- Effective coupling constant in every vertex is given by an integral of wave functions, e.g. for 3 V vertex:
$g_{(a, i)(b, j)(c, k)}=g_{5} \int_{0}^{\pi R} \varphi_{a, i} \varphi_{b, j} \varphi_{c, k} \mathrm{~d} y$
- Scattering amplitudes of general $V_{\mathrm{L}} V_{\mathrm{L}} \rightarrow V_{\mathrm{L}} V_{\mathrm{L}}$ process do not grow explicitly with energy for any allowed choice of boundary conditions.


## Example

Boundary conditions:

$$
\begin{aligned}
& \text { - } y=0: A_{\nu}^{1,2}=0, \partial_{y} A_{\nu}^{3}=0 \\
& \text { - } y=\pi R: \partial_{y} A_{\nu}^{a}=0
\end{aligned}
$$

Wave functions and masses:

- $\varphi_{W, n}(y)=\sqrt{\frac{2}{\pi R}} \sin \frac{(2 n-1) y}{2 R} \Rightarrow m_{W, n}=\frac{(2 n-1)}{2 R}, n \geq 1$
- $\varphi_{Z, n}(y)=\sqrt{\frac{2}{2^{\delta_{k, 0} \pi R}}} \cos \frac{2 n y}{2 R} \Rightarrow m_{Z, n}=\frac{2 n}{2 R}, n \geq 0$

Coupling constant:

- $g_{(W, i)(W, j)(Z, k)}=\frac{g_{5}}{\sqrt{2^{\delta_{k, 0} \pi R}}}\left(\delta_{j, i+k}+\delta_{j, i-k}-\delta_{j, 1-i+k}\right)$


## Final remarks

- There are several ways how to break electroweak symmetry in higher-dimensional models.
- 5D theory is non-renormalizable, unitarity breakdown is postponed to the cutoff scale of effective 4D theory (related to the size of extra dimension).

Realistic models:

- More complicated gauge group, e.g. $S U(2)_{R} \times S U(2)_{L} \times U(1)_{B-L}$
- Warped extra dimensions.


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## Thank you for your attention.

Vector bosons scattering and boundary conditions in KK toy model arXiv:1106.3043 [hep-th]

