

# Color Sextet Scalars in Colliders, Cosmology and low energy physics



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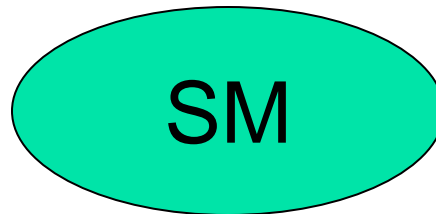
# Plan:

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- TeV mass color sextet scalars (*über-Higgses*)
- Neutrino mass + quark-lepton unification → a motivation for Color sextet scalars
- Can have TeV to sub-TeV mass;
- **Lead to new phenomena:**
  - (i) neutron-anti-neutron oscillation- no proton decay
  - (ii) new way to understand the origin of matter;
  - (iii) FCNC effects
  - (iv) New signals at LHC

# NEW TEV SCALE SCALARS

- **Bottom-up view:** What are possible TeV mass scalars that could couple to SM fermions (quarks) without making trouble ?



- Color quantum no.  $\{1\}$   $\{3, 3^*\}$   $\{6, 6^*\}$   
 couples to  $\bar{Q}u, \dots$   $QQ, QL$   $(QQ)_6$
- SM Higgs  $\otimes$  Leads to p-decay **Allowed; are they there ?**

# Their interaction and new class of baryon violation

- Three types of sextet scalars:

$$\Delta_{uu}, \Delta_{ud}, \Delta_{dd}$$

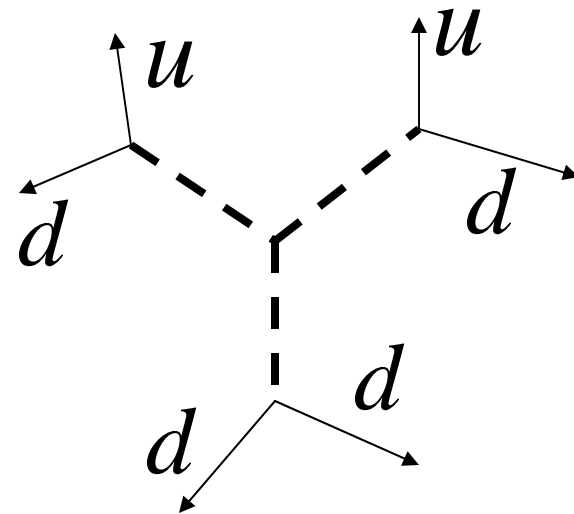
- Their interactions:

$$\mu \Delta_{ud} \Delta_{ud} \Delta_{dd} +$$

$$L = f_{uu} u^c u^c \Delta_{uu} + f_{ud} u^c d^c \Delta_{ud} + f_{dd} d^c d^c \Delta_{dd} + h.c.$$

- Leads to  $B=2$  baryon violation:  $n \leftrightarrow \bar{n}$

$$G_{B=2} \sim \frac{f^3 \mu}{M_{\Delta}^6}$$



# Sextet scalars from Seesaw Paradigm for neutrinos

- Seesaw mechanism: SM+ Majorana  $N$

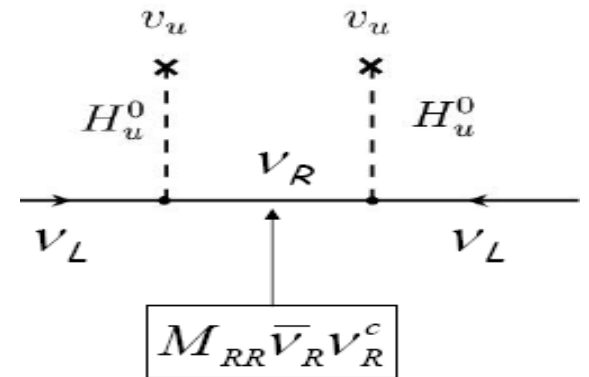
$$L_Y = h_\nu \bar{L} H N_R + M_R N N$$

- $M_R$  New scale

$$m_\nu \cong -\frac{h_\nu^2 v_{wk}^2}{M_R}$$

-Neutrino majorana

→ small nu mass natural since  $M_R \gg h v_{wk}$





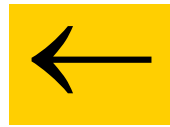
# Origin of scale $M_R$

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- Standard model: gauge sym.  $SU(2)_L \times U(1)_Y$

- Fermions:

$$m_\nu = 0$$



$$\begin{pmatrix} u_L \\ d_L \\ \nu_L \\ e_L \end{pmatrix} \quad \begin{matrix} u_R \\ d_R \\ e_R \end{matrix}$$

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- $N_R \rightarrow$  Gauge group:

$$SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$$

$$\begin{pmatrix} u_L \\ d_L \\ \nu_L \\ e_L \end{pmatrix} \xleftrightarrow{P} \begin{pmatrix} u_R \\ d_R \\ \nu_R \\ e_R \end{pmatrix}$$

- New scale - breaking

# New Scalars for symmetry breaking and fermion masses

$$\phi = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix} \quad \Delta = \begin{pmatrix} \frac{1}{\sqrt{2}} \Delta^+ & \Delta^{++} \\ \Delta^0 & -\frac{1}{\sqrt{2}} \Delta^+ \end{pmatrix}$$

$$\phi = \begin{pmatrix} \kappa & 0 \\ 0 & \kappa' \end{pmatrix} \quad \langle \Delta \rangle = \begin{pmatrix} 0 & 0 \\ v_R & 0 \end{pmatrix}$$

Break symmetry

$$L = h \bar{Q}_L \phi Q_R + f (LL\Delta_L + RR\Delta_R) + h.c.$$





# SEESAW FROM LR

$$SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$$

Nu-masses

$$\begin{pmatrix} 0 & 0 \\ \overline{0} & \overline{M} \end{pmatrix}$$



$$\langle \Delta_R \rangle \neq 0 \\ M_{W_R}, M_{Z'} \neq 0$$

$$SU(2)_L \otimes U(1)_Y$$



$$\langle \phi \rangle = \begin{pmatrix} \kappa & 0 \\ 0 & \kappa' \end{pmatrix}$$

$$\begin{pmatrix} 0 & m \\ \overline{m} & \overline{M} \end{pmatrix}$$

$$U(1)_{em}$$

$$M_{W_L}, M_Z \neq 0; m_{q,l} \neq 0$$

**Seesaw matrix:** seesaw  
scale is LR scale

# Quark-lepton Unification and color sextets

■ LR+QL  $\rightarrow$

$$\begin{pmatrix} u & u & u & \nu \\ d & d & d & e \end{pmatrix}_{L,R}$$

■ Gauge group:  $SU(2)_L \times SU(2)_R \times SU(4)_C$  (Pati, Salam'74)

■ **New Higgses for seesaw:** (Mohapatra, Marshak'80)

■  $\Delta_{ll} LL \rightarrow \Delta_{qq} QQ$  **Color sextets**

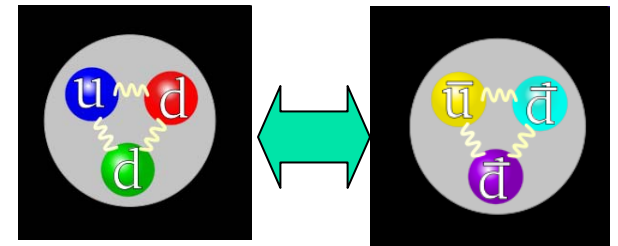
■  $(1, 3, 2) \rightarrow (1, 3, 10)$   
+ lepto-quarks

# New phenomena associated with sextets

- n-n-bar oscillation but no proton decay
- Majorana neutrino  $\rightarrow L=2$ ;
- n-n-bar oscillation  $\rightarrow B=2$  (Mohapatra, Marshak'80)

- n-n-bar diagram:

$$A_{n\bar{n}} \sim \frac{\lambda f_{qq}^3}{M_\Delta^5}$$

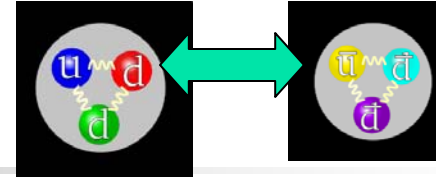


$$\delta_{n\bar{n}} = A_{n\bar{n}} \Lambda_{QCD}^6 \Rightarrow \tau_{n\bar{n}} = \frac{\hbar}{\delta_{n\bar{n}}}$$

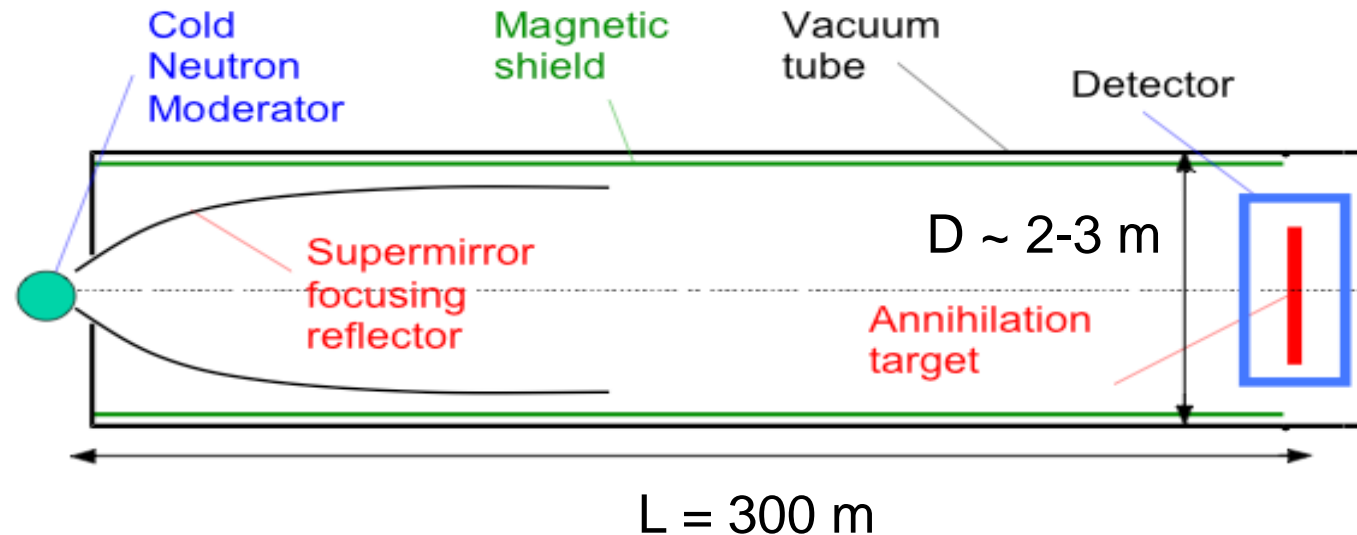
- Oscillation time:

- $M_\Delta \sim TeV$        $\tau_{n\bar{n}} \cong \text{yr to few years}$  ; **Observable!**

# Search for n-n-bar Osc. current status



- Free neutron oscillation in reactors: generic setup



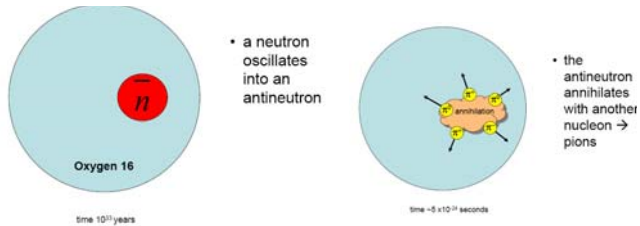
- Current bound (ILL'94)

with  $L \sim 90$  m and  $\langle t \rangle = 0.11$  sec  
measured  $P_{n\bar{n}} < 1.6 \times 10^{-18}$   
 $\tau > 8.6 \times 10^7$  sec

- Searching for new site: NANO workshop'Feb.'11

# Bound neutron:

- Searches in nucleon decay expts: S-K, Soudan, IMB...

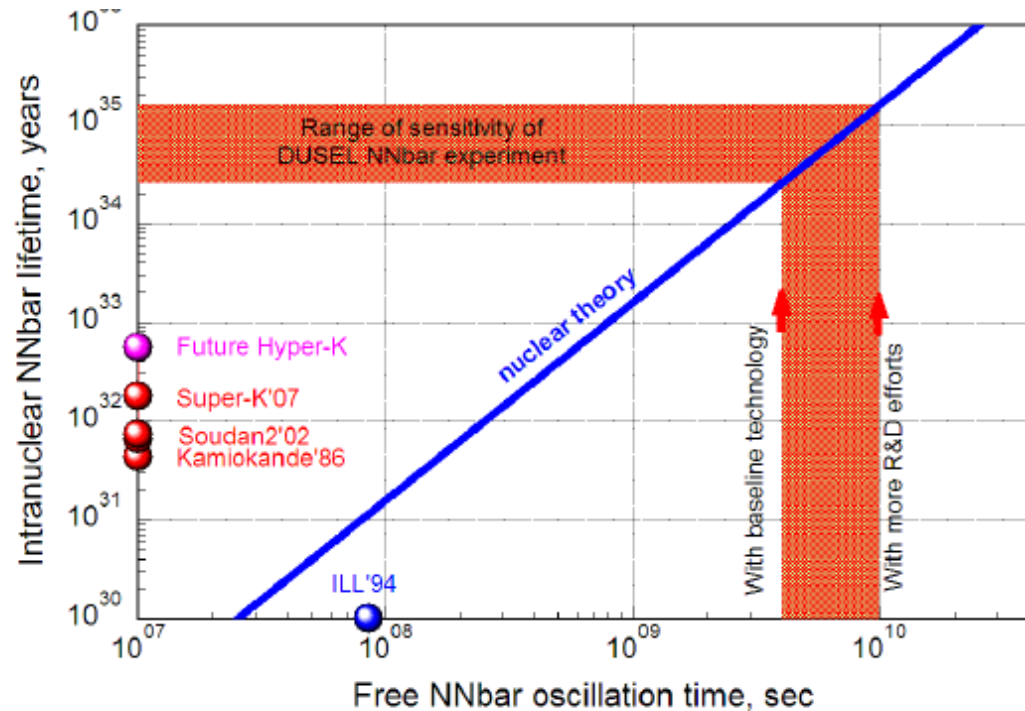


$$\tau_{Nuc} = R \tau_{free}^2$$

$$R = 0.3 \times 10^{23} \text{ sec}^{-1}$$

- Future

$\tau_{nucl} \rightarrow 2 \cdot 10^{32} \text{ yr}$  (SK-2009)

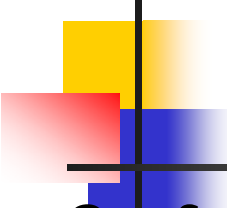




# n-n-bar as complementary to proton decay search

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- Proton decay will search B-L (Seesaw) scale in the range of GUT scale ( $\sim 10^{14}$  GeV or higher);
- Neutron-anti-neutron oscillation will search B-L from few TeV (without susy) to  $10^{11}$  GeV (with susy).
- Urge new direct search for NNbar.
- There is a proto-collaboration (**NANO**)



# Can diquarks be naturally light in high B-L scale models

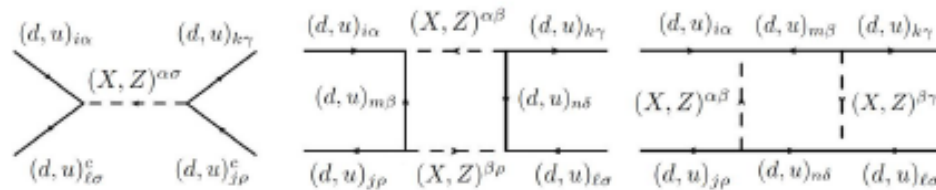
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- So far we have assumed that seesaw scale is also in the TeV range.
- In susy seesaw models, even for seesaw scale in the  $10^{10}$  range, due to accidental global symmetry of the superpotential,  $\Delta_{u^c u^c}$  can be naturally in TeV range

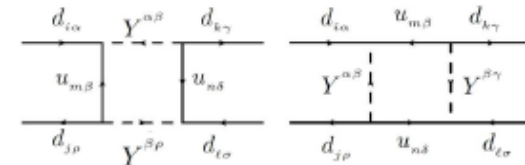
(Chacko, Mohapatra'98)

# Light diquarks and FCNC

- Tree and box diagrams lead to FCNC effects



(dd and uu sextets)



ud sextets:

- Tree graph constraints:** (Babu, Fortes, RNM'2011(to appear) )

--TeV mass dd-sextets:  $f_{11}^{dd} \neq 0$ ;  $f_{22,33}^{dd} \approx 0$  (at least 2 diags tiny)  
 (from B-B-bar and K-K-bar)

--TeV mass uu-sextets:  $f_{22}^{uu} \approx 0$ ;  $f_{11,33}^{uu} \neq 0$   
 (D-D-bar)



# Light diquarks-off-diagonal couplings

- Only box diagrams lead to FCNC effects:

- Allowed domain:

--TeV mass dd-sextets:  $|f_{i3}f_{i1}^*| < 0.0302$

(from B-B-bar and K-K-bar)

$$|f_{i2}f_{i1}^*| < 0.0147$$

--TeV mass uu-sextets:

(D-D-bar)

$$|f_{i2}f_{i1}^*| < 0.0111$$

- Many allowed textures for uu coupling: e.g.

$$f_{ij} = \begin{bmatrix} 0.3 & 0 & 0.3 \\ 0 & 0 & 0 \\ 0.3 & 0 & 0.3 \end{bmatrix}.$$

# Color sextets and some anomalies:

- Sextet induced Bs-Bs-bar mixing with CP violation in its coupling can explain Tevatron dimuon anomaly:
- There is also an anomaly in understanding  $\varepsilon_K$  in SM (Buras-Guadagnoli) :  $\varepsilon_K$  and  $B \rightarrow J/\psi + K_S$  involve  $\sin 2\beta = 0.79 \pm 0.039$   $0.672 \pm 0.023$   
Color sextet contribution can resolve this.
- Anomaly in B-decay-(sextets can resolve)  
 $A_{CP}(B^0 \rightarrow K^+ \pi^-) = -9.7 \pm 1.2\%$ ;  $A_{CP}(B^+ \rightarrow K^+ \pi^0) = 5.0 \pm 2.5\%$
- Explaining decay anomalies require ud mass < TeV

# Color sextet scalars at LHC

- Two production modes at LHC:

(I) **Single production**  $uu \rightarrow \overline{\Delta}_{u^c u^c} \rightarrow tt$  or  $t + \text{jet}$

xsection calculated in (RNM, Okada, Yu'07;) resonance peaks above SM background- decay to  $tt$  or  $tj$   
**depending on color sextet coupling texture;**

- CDF limit with  $6.1 \text{ fb}^{-1}$   $\sigma(tt + \bar{t}\bar{t}) \times \text{Br}(W \rightarrow l\nu)^2 < 54 \text{ fb}$ ,

(II) **Drell-Yan pair production:**  $q\bar{q} \rightarrow G \rightarrow \Delta_{u^c u^c} \Delta_{u^c u^c}^*$

**model indep.** (Chen, Klem, Rentala, Wang, 08)

- Leads to  $t\bar{t}\bar{t}$  final states: **LHC reach < TeV**

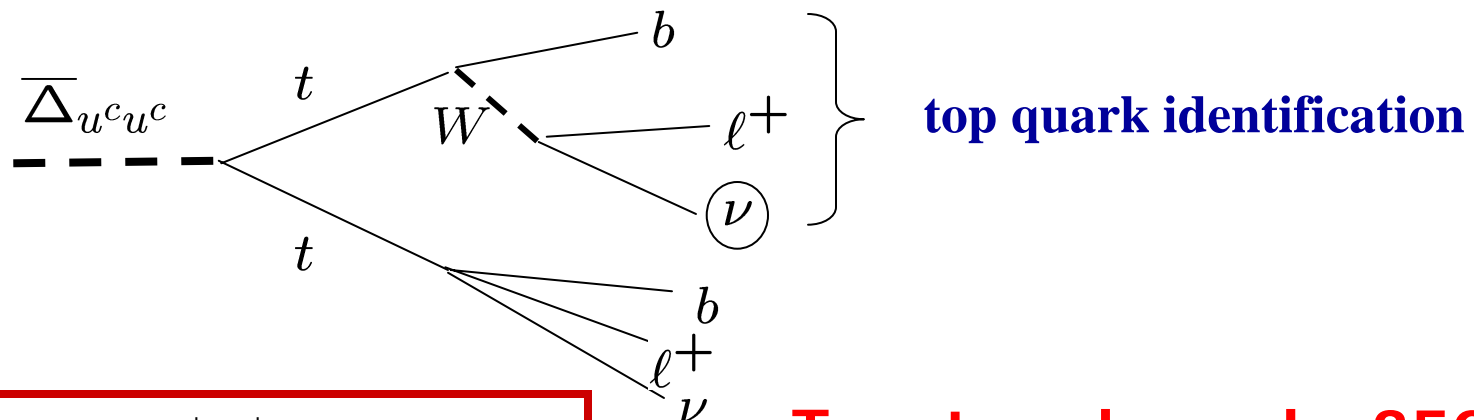
(Berger, Cao, Chen, Shaughnessy, Zhang'10; Han, Lewis'09)

# Signal for single production

So, our target is

$$\begin{cases} uu \rightarrow \overline{\Delta}_{ucuc} \rightarrow tt \text{ or } t + \text{jet} \\ \bar{u}\bar{u} \rightarrow \overline{\Delta}_{ucuc}^* \rightarrow \bar{t}\bar{t} \text{ or } \bar{t} + \text{jet} \end{cases}$$

To measure diquark mass (final state invariant mass)



$$pp \rightarrow l^\pm l^\pm bb + Nj + ..$$

- Tevatron bound  $> 350$  GeV.
- LHC can probe till TeV.

# SINGLE SEXTET PRODUCTION AT LHC: (MODEL DEP.)

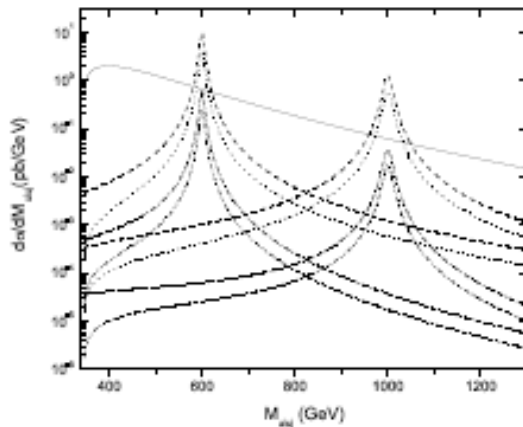
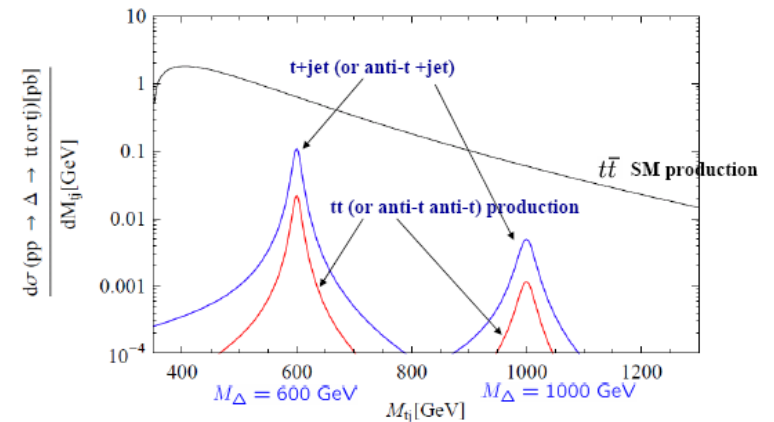


FIG. 2: The differential cross sections for  $tj$  (dashed line),  $tt$  (dotted line),  $\bar{t}j$  (dashed-dotted line) and  $\bar{t}\bar{t}$  (dashed-dotted-dotted line) as a function of the invariant mass of final state  $M_{u_i u_j}$ . The left peak corresponds to  $m_\Delta = 600$  (GeV) and the right one to  $m_\Delta = 1$  TeV. The solid line is the standard model  $t\bar{t}$  background.

$E_{CMS} = 14$  TeV



Production from sea quarks

**Diquark has a baryon number & LHC is ``pp'' machine**

$$\rightarrow \sigma(tt) \gg \sigma(\bar{t}\bar{t}), \quad \sigma(t + jet) \gg \sigma(\bar{t} + jet)$$

**Depends on Yukawa coupling**

$$f_{ij} = \begin{bmatrix} 0.3 & 0 & 0.3 \\ 0 & 0 & 0 \\ 0.3 & 0 & 0.3 \end{bmatrix}.$$

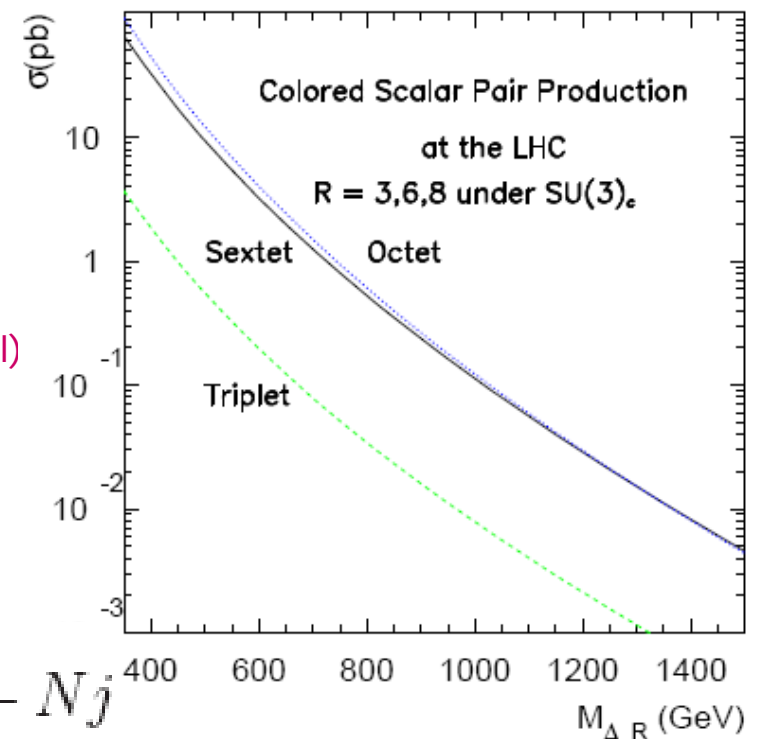
# Pair Production of Deltas


- Due to color sextet nature, Drell-Yan production enhanced- independent of Yukawa coupling  
(model indep)

- Leads to  $t\bar{t}t\bar{t}$  final states (Chen et al)

- Signal

$$pp \rightarrow \bar{\Phi}_6 \Phi_6 \rightarrow t\bar{t}t\bar{t} \rightarrow 4b + e^\pm e^\pm + \cancel{E}_T + Nj$$





# Observable $n\bar{n}$ and baryogenesis

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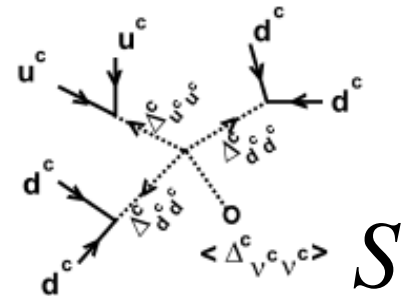
- If  $n\bar{n}$  is observable, it erases any preexisting baryon asymmetry above 100 GeV.
- New way to understand must involve  $N\bar{N}$ -bar necessarily.
- Note  $N\bar{N}$  being a high dim operator goes out of eq. below EW scale- must involve baryogenesis without help from sphalerons.

**Post sphaleron baryogenesis**

# New way to understand Origin of matter

- 224 model realization of color sextets: there is a neutral field accompanying color sextets,  $S$  with effective int.

$$\frac{S}{\Lambda^6} u_R d_R d_R u_R d_R d_R$$



$S$  acquires vacuum expectation value and breaks Baryon number

$S_r$  decays to 6 quarks as well as into 6 antiquarks  $S_r \rightarrow 6q$ ,  $S_r \rightarrow 6\bar{q}$

Leads to baryon asymmetry

Babu, Nasri, RNM'06; Babu, Dev, RNM'09





# Post sphaleron baryogenesis

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- Due to high dimension, S-decay occurs below EW scale; → **after sphaleron decoupling**.
- Requires light sextets ( $\sim$ TeV mass)
- For a certain parameter domain of the model, adequate baryon asymmetry implies  $n\bar{n}$ -bar transition time  $< 10^{10}$  sec. (**observable using current reactor fluxes**)



# Conclusion:

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- **Color sextets** in the TeV range are motivated by seesaw paradigm for neutrino masses and quark-lepton unification (hint from  $m_b \cong m_\tau$ ;  $m_\mu = 3m_s$  at M)
- **Lead to observable NNbar oscillation and new mechanism for baryogenesis.**
- **LHC-an ideal machine to search for sub-TeV mass.**
- **nn-bar** complementary to proton decay for understanding the seesaw physics and new direct search important for true understanding of seesaw and B-L symmetry.